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1971

# A study of the preparation, structure and properties of compounds containing polynuclear cations of bismuth: an investigation of anion effects

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#### 71-21,942

FRIEDMAN, Robert Mark, 1945- A STUDY OF THE PREPARATION, STRUCTURE, AND PROPERTIES OF COMPOUNDS CONTAINING POLYNUCLEAR CATIONS OF BISMUTH. AN INVESTIGATION OF ANION EFFECTS.

Iowa State University, Ph.D., 1971 Chemistry, inorganic

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A study of the preparation, structure, and properties. of compounds containing polynuclear cations of bismuth.  $\bullet$ An investigation of anion effects

by

Robert Mark Friedman

**A** Dissertation Submitted to the Graduate Faculty in Partial Fulfillment of The Requirements for the Degree of DOCTOR OF PHILOSOPHY

Major Subject: Inorganic Chemistry

Approved:

Signature was redacted for privacy.

In Charge of Major Work

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#### INTRODUCTION

Following World War Two science was deluged with financial support because of the realization of the impact that fundamental research can have on technology. The study of molten salts was the recipient of much attention owing to their unique physical properties, e.g., high thermal stability, low vapor pressure, wide range of accessible temperatures, good electrical conductivity, and low viscosity.

Interest in the use of molten salts in nuclear reactors as solvents for fissionable materials or as a heat exchange medium, industrial concern about the mechanism of corrosion, and the low current yields in electrometallurgical processes provided the impetus for basic investigations into the nature of the solution of metals in their molten salts. Several theories, differing mainly in the character of the solute species, arose to explain the phenomenon. The three most popular solution models were: A.) the dissolution of metal atoms such as found in the vapor phase, B.) the solution of "ion pairs" with solvated electrons occupying anion sites in the quasilattice of the melt, and C.) the formation of a reduced solute cation by an oxidation-reduction reaction between the metal and the melt cations. None of these hypotheses alone was capable of explaining all of the observed data. It was sometimes necessary to choose a combination of models to explain all of the

physlochemical properties of a single system, i.e., a different model for a different property.

One of the more intricate systems to be Investigated was the bismuth-bismuth trichloride system. It had long been known (1) that when bismuth dissolved in its trichloride a black solid was formed with the approximate composition BiCl. Although simple monomeric BiCl was observed and identified (2) in the gas phase in 1931, It was found necessary to Invoke polymerization to explain the dlamagnetlsm of the solid isolated by Corbett (3). The properties of dilute solutions of the metal in the molten halide were intensively studied. The liquidus of the BiCl-BiCl<sub>3</sub> portion of the phase diagram  $(4)$ was interpreted  $(3)$  in terms of a  $(BiCl)_n$ , solute, whereas precise cryoscoplc measurements (5) suggested the presence of  $(B1C1)_2$  or  $B1\frac{2}{3}$ . Vapor pressures (6) over solutions containing up to 24 mole % metal were accounted for by considering alternatively metal atoms  $(6)$ , tetramers of the form  $(B1C1)$ ,  $(7)$ , or a mixture of  $(BICl)_2$  and  $Bl_2$  (8) to be present. A reexamination of the phase diagram of the bismuth-bismuth trichloride system (9) essentially confirmed the earlier work by Sokolova (4) on the liquid-solid region. The metal and salt were found to be completely miscible above  $780^\circ$ .<sup>1</sup> In

<sup>1</sup>All temperatures are in degrees Centigrade unless otherwise noted.

determining the features of the two liquid region, the occurrence of retrograde solubility of bismuth on the salt-rich side of the miscibility gap was discovered; the solubility of bismuth in its trichloride decreased from  $45\%$  at 320° to 28% at about 550®. The initial decrease in solubility followed by its Increase as the temperature was increased to the consolute temperature was interpreted as the resultant of two processes, one predominant at lower temperatures and the other predominant at higher temperatures; for example, the solution of bismuth as a polymeric subhallde at lower temperatures followed by decomposition at higher temperatures. Polarographlc measurements (10) on very dilute solutions gave results consistent with a  $B1^+$  solute species, while EMF concentration cells (11) showed  $B1^+$  to be present below 0.6 mole % metal with polymerization to a species such as  $B1_4^{1+}$  at higher metal concentrations, A spectrophotometrlc study (12) established the presence of two chromophores in the range 0.01 to 6 mole  $%$ metal. An unambiguous identification of the solute species was not possible, but the simplest equilibrium consistent with the data (13),  $4B1^+$   $\stackrel{+}{\leftarrow}$  B1<sup>1</sup><sup>+</sup>, was in agreement with the EMF measurements.

Corbett and McMullan (14) prepared BiAlCl4, analogous to BlCl, by adding aluminum chloride to the metal-metal hallde binary system. They found complete conversion to the "acidstabilized" salt at 260° In contrast with the 46% conversion to BiCl in the binary. The compound melted at  $253^\circ$ . It was

maroon In bulk and reddish-brown as a powder, but darkened rapidly in air. Ah X-ray diffraction study (15) indicated the probable space group as R3c for BiAlCl4. The radial pair distribution functions of the melt and of the powder of this composition were interpreted in terms of an equilateral triangle of bismuth atoms separated by  $3.04$  Å (16). Based on the assumed stoichlometry the Investigators concluded that the correct formulation of the compound should be  $(Bi<sub>3</sub>)<sup>3+</sup>(A1C1<sub>4</sub>)<sub>3</sub>$ . The probable existence of polynuclear cations of bismuth appeared to be established both in the melt and the solid state.

The unravelling of the crystal structure of bismuth monochlorlde (17) revealed a greater breadth and complexity in the chemistry of bismuth. The existence of a discrete  $B1\frac{5}{9}$ <sup>+</sup> ion with two types of chlorobismuthate(III) anions resulted in an overall stoichiometry of BiCli.167. A rationalization of this strange beast was proposed by Corbett and Rundle **(I8)** through an extended Huckel LCAO-MO model utilizing a 6p atomic orbital basis set on each bismuth. The success of this simple model to provide a reasonable description of the electronic structure of this unusual cation encouraged its application to similar problems.

Previous workers had proposed the existence of  $B1<sup>3</sup>$  ions in solid and molten BiAlCl<sub>4</sub>. Corbett (19) noted that this ion did not have a reasonable bonding scheme for a dlamagnetic compound since molecular orbitals with symmetries  $a_1'$ ,  $a_2''$ , and  $e'$  were all strongly bonding; thus,  $B1\frac{1}{3}$ , with eight bonding

electrons, would be predicted. This Ion would also be compatible with the X-ray diffraction results as well as studies of dilute solutions of bismuth in bismuth trichloride. In the dilute solution studies one could not distinguish between the family of solutes  $(Bi_3^{\dagger} \cdot nBi^{\dagger 3})$ ; hence,  $Bi_3^{\dagger}$  and  $Bi_4^{\dagger +}$ , previously proposed to fit the data, would be equivalent.

Investigation of the  $B1-B1Cl<sub>3</sub>$  -3AlCl<sub>3</sub> system uncovered two new compounds,  $Bi_5(A1CL_h)$ , and  $Bi_4(A1CL_h)$ , and a glass phase which remained uncharacterized (19). Concurrently a spectrophotometrlc study of the dilute solutions of bismuth and bismuth trichloride in NaCl-AlCl<sub>3</sub> and KCl-ZnCl<sub>2</sub> solvents revealed the presence of  $B_1^+, B_1^3^+,$  and  $B_1^2^+$  ions (20,21). Corbett noted the spectral resemblance between the above solid compounds and the two polymeric ions present in the melt and concluded that they were present in the solid phases. Using the simple Huckel molecular orbital procedure and the observed dlamagnetlsm of the compounds he predicted the trigonal bipyramid,  $D_{3h}$  symmetry, and the Archimedean antiprism,  $D_{4d}$ symmetry, to be the geometries of the  $B15^+$  and the  $B18^+$  ions, respectively. The other reasonable geometries for five and eight atom polyhedra had open-shell electronic configurations. The powder pattern for  $B1_5(A1Cl_4)$ <sub>3</sub> accounted for the reported diffraction pattern of "BiAlCl<sub>4</sub>" and the triangular units deduced earlier were consistent with the proposed geometry. All the faces of the regular trigonal blpyramld are triangular and the other distance between atoms in that polyhedron,

between the apices, was probably too long to be resolved by the radial distribution function. The "BiAlCl<sub>4</sub>" composition is close to the liquidus composition at the temperature used to initially prepare it.

The  $B1^+$  ion has also been observed spectrally (22) in the NaBr-AlBr<sub>3</sub> eutectic. The spectrum of  $Bi^+$  in both the chloride and bromide media has been successfully rationalized (23) by a ligand field treatment of the perturbation of the  $6p^2$  atomic states in an environment of lower than cubic symmetry. The relativistic treatment predicted a  ${}^{3}P_{0}$  diamagnetic ground state would be well separated from the excited states.

The discovery of the  $B1_5^{3+}$  and  $B1_6^{2+}$  ions demonstrated that the original polynuclear cation of bismuth,  $B_1^{\frac{1}{9}}$ , was not unique in specie, but part of a general class of heretofore unrealized catlonic species. It was noticed that each of the two systems studied had uncovered different cations. The intent of the work reported herein was to be the investigation of factors affecting the stability of the polynuclear cations with respect to different anion environments. This was to be achieved by using different Lewis acids than aluminum chloride to prepare them. Changing the Lewis acid present in the melt changes the pCl", hence the polarization and coulombic interactions in the synthesis media, and, perhaps as importantly, changes the available anions for formation of an isolable solid containing these species.

To consider the possible variations of anions and their effects on the polynuclear cations, the electrostatic model was chosen as an appropriate framework; the bismuth cations were viewed as being discrete in the solids. The leading terms in this analysis are the coulombic and repulsive interactions represented by the "hard sphere" ionic model. Other terms, usually considered less important in cases involving mononuclear ions, represent dipole, induced dlpole, and multipole interactions. The polarizablllty of the ions were viewed as being indicative of the magnitude of these last effects.

In a discussion of the stability of reduced metal halldes to disproportionation, Johnson (24) has shown, in the hard sphere ionic limit and within the validity of the Kapustlnskll equation, that the stability of a mononuclear lower oxidation state of a metal increases with Increasing size of an anion of the same charge. For the common lattice types the Kapustlnskll equation, through Judicious reparameterization, eliminated the structural dependence of lattice energies from their calculation. In addition to noting the generalization reached by Johnson, the manner in which the particular geometry of the anion influences the ionic packing in the solid was considered another variable in this study.

The stabilization of lower oxidation states in solids using Lewis acid complexatlon was first demonstrated for  $Cd_2(ALCl_*)_2$  (25). Further work has indicated that zinc dichlorlde and beryllium dichloride do not enhance the

reduction of cadmium dichlorlde to the monovalent state (J. D. Corbett and W. J. Burkhard, 1968, personal communication). These results discouraged the use of these salts in the present investigation. The  $KCI-ZnCl<sub>2</sub>$  solvent used in the spectral identification of the  $B_1^+$  and  $B_1^3$ <sup>+</sup> ions contained 72 mole %  $ZnCl<sub>2</sub>$ , and therefore the  $ZnCl<sub>3</sub>$  and  $Zn<sub>2</sub>Cl<sub>5</sub>$  ions were probably the predominant anions present in the melts.

Hafnium tetrachloride and tantalum pentachloride were chosen as the Lewis acids for initial study because of their availability, the stability co reduction of each metal in its highest oxidation state (26,27), and the presumed existence of only the hexachlorometallate ions in fused chloride media (28). The possible comparisons of the effects of charge, size, and geometry between the anions HfCl<sup>2</sup>; TaCl<sub>6</sub>, AlCl<sub>4</sub>, BiCl<sup>2</sup><sup>-</sup>, and  $B1_2C1_6^2$  were considered maximal. Other research was later conceived to clarify the results obtained in these systems.

#### EXPERIMENTAL PROCEDURES

#### Materials

The bismuth metal used was American Smelting and Refining 99.999% grade. Spectrographlc analysis showed the presence of Cu and Ag at the trace level, and Pb, Pe, and Si at the faint trace level. The metal was vacuum fused to allow decomposition and physical removal of impurities, and then rebroken into small pieces for ease in handling and weighing. The metal was stored under vacuum; however, transfers and weighings were performed In the air.

Crystal bar hafnium was supplied by the Pittsburg Naval Reactors Office, Emission analysis revealed the metal to contain 1.4 weight per cent Zr, with Cu and Fe as trace Impurities, and Mg and SI as faint trace impurities.

Two mil tantalum sheet was obtained from Fansteel with typical analyses showing the following impurity levels in p.p.m.; 0, 15; N, 45; H, 5-; W, 25-; C, Nb, Zr, Mo, Tl, Fe, Ni, SI, Mn, Ca, Al, Cu, Sn, Cr, V, Co, Mg, 10-.

Commercial bismuth trichloride was purified by dehydration at 100°, distillation under a partial pressure of oxygen, and subsequently sublimation under dynamic vacuum. This procedure was employed to remove the carbon impurity present in the commercially available material. An attempt was made to omit the distillation under oxygen to avoid the excessive formation of bismuth oxychloride; however, even

repeated sublimations through a course-grade slntered-glass disk were not as successful in removing the carbon impurity. Approximately a 33% yield of white crystalline bismuth trichloride was obtained through this process. In spite of the low yield, this procedure was considered practical in terms of effort and cost because the alternative method, direct reaction between the elements, was known to be extremely slow even under forcing conditions.

Dry box and vacuum line techniques were used throughout this investigation because of the sensitivity of the salts to oxygen and moisture. The drybox was continuously flushed with argon and its atmosphere recycled through Linde molecular sieves. To further Insure a low partial pressure of water, an open tray of phosphorus pentoxide deslccant was kept inside the dry box. An evacuable port was used to facilitate entry. All manipulations and weighings of salts were performed in the dry box. The weighings were made using a triple beam, single pan balance which was accurate to five milligrams. Salts were stored in evacuable containers or sealed evacuated glass ampoules.

Commercial potassium chloride, dried under dynamic vacuum at 300°, was used without further purification. Other commercial salts, aluminum tribromide, bismuth tribromide, and ammonium chloride, were sublimed twice under dynamic vacuum before use. The pure salts were white, bright yellow, and white, respectively.

Chlorine gas was obtained in lecture bottles from Mathegon Co. Electronic grade hydrogen chloride was supplied by Precision Gas Products.

#### Synthesis

Hafnium tetrachloride was prepared from hafnium metal and hydrogen chloride gas In the apparatus pictured in Figure la. The main chamber of the apparatus was constructed with the hafnium metal contained Inside it. The apparatus was flushed with a stream of HCl to purge it of air and moisture which would react with the product. The metal was slowly heated to **360°** while the gas continued to pass over it. The hafnium tetrachloride formed by the reaction at point B collected in the cooler region of the compartment at site C. When the reaction was complete the stopcocks at A and D were closed and the tube was evacuated on a vacuum line. The container was taken into the dry box where It was cracked open and the salt was transferred to a sublimation vessel. The salt was sublimed under dynamic vacuum and then stored.

Tantalum pentachlorlde was prepared from the elements. Strips of tantalum were loaded into compartment C of the apparatus shown In Figure lb. Chlorine was distilled into compartment A which was cooled using a Dry Ice-acetone slush bath. The apparatus was sealed off and placed in a furnace which encased the central chamber and parts of the neighboring chambers (B and D). Chlorine was moved back and forth between

## Figure 1. Apparatus for the preparation of (a) hafnium tetrachloride and (b) tantalum pentachloride

 $\mathcal{L}^{\text{max}}_{\text{max}}$  and  $\mathcal{L}^{\text{max}}_{\text{max}}$  . The  $\mathcal{L}^{\text{max}}_{\text{max}}$ 

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 $\sim 10^{11}$  km  $^{-1}$ 

 $\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}})$  and  $\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}})$  . The contribution



the cold traps (A and E) passing over the metal heated first at 375° and later at 425°. The pentachlorlde which collected in the cooler regions of the tube had a yellow cast whereas the pure material Is reported to be white. Sublimation of the product yielded crystals which retained a slight yellow tinge. This coloration may have been due to the presence of small amounts of yellow WCl<sub>5</sub> introduced by chlorination of the tungsten Impurity present In the tantalum sheet metal.

Reduced bismuth phases were synthesized by the reduction by elemental bismuth of mixtures of bismuth trichloride and the appropriate metal chloride, HfCl, or TaCl<sub>s</sub>, The composition of the unreduced mixture was determined by the requirement that all chloride be formally and stolchlometrically converted to the larger less basic MCl $_{6}^{n+}$  anion present in the respective system. Previous investigations have shown, however, that there are no intermediate compounds formed in either the BiCl<sub>3</sub>-HfCl<sub>i</sub> (29) or the BiCl<sub>3</sub>-TaCl<sub>5</sub> (30) systems.

Reaction containers were constructed of fused silica to prevent side reaction between the melts and the container walls. Either 15 or 40 mm. dia. tubing was closed on one end and connected through 8 mm. tubing to a ball joint on the other end. The choice of the diameter of the reaction vessels was made according to the desire for either a low volume and less material in the vapor phase or a large interface between the metal and salt phases. When separation of crystals from the melt was necessary, a two compartment 20 mm. dia. tube was

used. Separating the two chambers was a silica disk with 2 mm. slits cut along a l80 degree arc. These tubes were loaded in the dry box and then sealed off under vacuum; thus, the reactions occurred under the vapor composition and pressure determined by the reactants. In order to minimize the thermal gradients the reaction tubes were Jacketed in a larger diameter glass tube and held in place using Pibrefrax (Carborundum Co.) insulation. The temperature was monitored by a thermocouple placed inside the glass Jacket near the sample. Cylindrical resistance furnaces used to heat the samples were regulated by Minneapolis-Honeywell Brown Electronlk Indicating Proportionating Controllers with a reported accuracy of  $\pm$  1°.

Samples of  $BI_{12}CI_{14}$  were synthesized by bismuth reduction of potassium chlorobismuthate melts. The reactants were heated to 340° to facilitate reduction and then equilibrated below 323°, the compound's incongruent melting point, for at least eight hours. The reduced melt was cooled at a maximum rate of five degrees per hour to 285° and then the reaction vessel was inverted filtering off the mother liquor.

Attempts to prepare  $B1_5(A1Br_4)_3$  were conducted under the assumption that the phase relationships in the Bi-BiBr3 \*3AlBr3 system were similar to those reported (19) for the analogous chloride system. To determine the appropriate temperatures for crystal growth, reactions and recrystallizations were carried out in a glass furnace. The furnace was made by winding 1 mm. Nichrome wire around a 30 mm. Vycor tube which

was Jacketed by a larger tube and held In place with asbestos tape. The advantage of being able to see the contents of the furnace was partly offset by the thermal gradients arising from the non-uniform spacing of the heating coil.

#### Analyses

Small glass vials were filled with samples and then capped while in the dry box. These sample containers were weighed in the air using an analytical balance accurate to 0.1 mg. Their contents were emptied into dilute nitric acid for dissolution and the empty vial weighed. The sample weights were determined by difference and were corrected for the buoyancy of argon. This correction was the difference In weights of a duplicate vial containing argon and then air.

Analysis for bismuth was carried out according to the method of Fritz (31) using EDTA with thiourea as the Indicator. Allquots were heated with concentrated perchloric acid until fuming to eliminate the chloride ions which prevent a sharp endpolnt. The pH was monitored throughout the titration because of the sensitivity of the bismuth-thiourea complex to the level of acidity.

Chloride was determined gravimetrlcally as AgCl from a 50:50 water-acetone solution.

After several attempts it was decided that conventional gravimetric analysis for potassium as the tetraphenylborate salt could not meet the desired accuracy. Potassium was then

determined by Mr. R, K. Hansen of the Ames Laboratory Analytical Services Group by atomic absorption spectroscopy using the 7664.9 Â resonance line.

#### Physical Measurements

The ultraviolet-visible spectra were recorded using a Cary 14 spectrophotometer. The samples were mulled with Nujol or petroleum Jelly and placed between two thin optical grade quartz plates. A thin bead of Dow Corning high vacuum silicone grease was placed around the edges of the plates to provide a seal against oxygen and moisture.

The far-Infrared spectra were measured on a Beckman Model IR-11 by Mr. C. Hill of the Ames Laboratory Spectroscopic Services Group. The samples were submitted as mulls placed between two sheets of polyethylene and secured in a cell equipped with an 0-rlng seal.

Magnetic susceptibility data were obtained over the temperature range 77°-300°K using the Faraday balance described in detail by Converse (32). The apparatus consisted of a Cahn model RQ electrobalance housed In an evacuable chamber and a Varian electromagnet. Samples were loaded into a small Teflon bucket with a screw-in cap. Measurements were made at five field strengths and susceptibilities corrected for ferromagnetic impurities were obtained using the Hondo-Owens equation. The data were processed on an IBM **360/65** computer using a Fortran IV program written by Converse. The treatment of the

data Included a dlamagnetlc correction for the sample container.

X-ray powder patterns were obtained using a Phillips Debye-Scherrer camera with a diameter of 114.59 mm. and exposure to nickel-filtered copper  $K_{\alpha}$  radiation. Powdered samples were loaded into 0.3 mm. Lindemann glass capillaries which were plugged with Aplezon Q wax in the dry box, and later sealed with a gas-oxygen hand torch. The d-spaclngs were read directly from the film with the aid of a calibrated template. Consideration was given to film shrinkage in the reported values. Intensities were estimated visually based on an arbitrary scale with a maximum value of 100,

#### Thermal Analysis

Sample containers for thermal analysis were constructed from 30 mm. lengths of 15 mm. o.d, fused silica tubing. The containers had a closed, 5 mm. length of 2 mm. tubing sealed inwardly through the bottom for a thermocouple well and a tiall Joint connected to the top through 8 mm. tubing. The containers were filled with three to eight grams of sample, sealed under vacuum, and Jacketed in another tube with the thermocouple in its well. The Junction potential of the thermocouple was plotted as a function of time on a Bristol Model **560** Dynamaster strip chart recorder and measured using a Rubicon potentiometer.

## Crystallographic Procedures<sup> $\perp$ </sup>

#### Crystal mounting

The crystals were taken Into a glove box specifically designed (35) for working under an inert atmosphere while viewing the work area through a long focal-length microscope. A Bausch & Lomb "Stereozoom" microscope with 0.5X objective lens and lOX eyepieces containing a calibrated scale was used. A removable Plexiglass top to the box allowed entrance to the interior and which, when secured, allowed evacuation of the chamber.

The samples were broken into small pieces and crystals carefully separated from the matrix. Crystals were cut into desirable size, approximately 0.1 mm. cubed, using a surgical scalpel. Before being cut the crystals were lightly coated with petroleum jelly to prevent the cut pieces from scattering. Suitable crystals were picked up on the end of a thin glass stalk and carefully inserted into 0.3 mm. Lindemann capillaries. The latter were held in a piece of 2 mm. l.d. Pyrex tubing which enclosed all of the fragile capillary except for the funnel-shaped open end. The glass stalk was removed from

 $^{1}$ Contrary to usual practice no attempt has been made to discuss the theory or practical aspects of crystallography beyond the material directly relevant to the problem covered by this thesis. The reader desiring further exposition is referred to the excellent texts by Buerger (33) and Stout and Jensen  $(34)$ .

the capillary leaving the crystal adhering to its walls. The open end of the capillary was then sealed using a resistance heated wire. When enough crystals had been mounted the capillaries were removed from the dry box, shortened using a gasoxygen hand torch, and their ends coated with Apiezon W wax as an additional protection against leaks. The capillary was held in a short piece of metal tubing using Apiezon W. This tube was then mounted in a goniometer head for X-ray measurements,

## $(B1^+) (B1\$ <sup>+</sup> $) (HfCl_6^2)$ ,

Preliminary Weissenberg and precession photographs taken with Ni-filtered Cu  $K_{\alpha}$  radiation showed the crystal system to be hexagonal and to belong to the 6/m Laue class. Extinction conditions,  $\lambda = 2n+1$  for 002 reflections, were consistent with space groups  $P6_3/m$  ( $C_{6h}^2$ , No. 176) and P6<sub>3</sub> ( $C_{6}^6$ , No. 173).

Precise lattice parameters and their standard deviations,  $a = 13.890 \pm 0.001$  and  $c = 10.692 \pm 0.002$  Å, were determined by least-squares fitting<sup>1</sup> (36) to the two theta angles (Mo  $K_{\alpha}$ radiation) of 21 Independent reflections. The diffraction angles were measured as the difference in the readings of the w arc of a previously aligned four-circle dlffractometer for

The computer programs used in this investigation are listed and gratefully acknowledged In Table 1. They were run on an IBM 36O Model **65** computer.



^Hubbard, C. R., Department of Chemistry, Iowa State University, Ames, Iowa. Private communication. **1969.** 

b<sub>Hubbard</sub>, C. R., and Quicksall, C. O., Iowa State University, Ames, Iowa. Private communication. I969.

Table 1. (Continued)

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^Hackert, M. L., Department of Chemistry, Iowa State University, Ames, Iowa. Private communication. 1969.

 $\frac{4}{3}$ 

the peak positions of hk& and hk& reflections. The center of each peak was determined by left-right, top-bottom beam splitting. By measuring the angle between the two positions of the crystal not only was the zero instrument error eliminated but also the errors due to absorption and incorrect centering of the crystal.

The density was measured micropycnometrically on a  $1 \text{ g}$ . sample by displacing chloroform previously dried over calcium hydride. The experimental value of  $6.05 \pm 0.2$  g/cm<sup>3</sup> compares favorably with a value of  $6.07 \text{ g/cm}^3$  calculated using  $Z = 2$ .

Three-dimensional peak height intensity data were collected on an automated Hilger-Watts four-circle diffractometer using Zr-filtered Mo  $K_{\alpha}$  radiation and a take-off angle of 4.5°. The crystal was mounted with the c-axis colinear with the spindle axis. The diffractometer and experimental arrangement have been described in detail elsewhere (41). Stationary crystal-stationary counter background measurements were made on both sides of the peak along a theta-two theta path. Periodic measurement of three standard reflections confirmed the absence of crystal decomposition or motion. A total of 1122 reflections were measured in the unique sector of a sphere bounded by  $\theta = 25^\circ$ . Approximately 100 reflections with diffraction angles spanning the range from zero to twenty-five degrees were remeasured using a G-20 step scan technique. These integrated intensities were used to convert the remaining reflections from peak height to integrated

intensities by the method of Alexander and Smith (42),

The net converted Intensities were also corrected for Lorentz, polarization, and absorption effects. The crystal used to gather data was a 0.1 mm. length of a needle with a distorted hexagonal cross-section, approximately 0.1 x 0.15 mm. Examination of the crystal under 60X magnification gave approximate dimensions. The precise shape and orientation of the crystal used to calculate the absorption correction (43) were varied within reasonable limits to minimize the difference between symmetry equivalent {hO&} and {Oh&} data. The transmission factors, computed using the program ABCOR (37) and a linear absorption coefficient of 571.0 ranged from **0.55%** to 4.33%, almost an eight-fold range.

While treating the raw intensities the standard deviation in each measurement was estimated as

 $\sigma_{\text{I}} = [C_T + C_B + (K_T C_T)^2 + (K_B C_B)^2 + (K_A C_R)^2]^{1/2}/A$ where  $C_{\phi}$ ,  $C_{R}$ ,  $C_{R}$  are the total, background, and net counts, respectively, and A is the transmission factor.  $K_{\text{m}}$ ,  $K_{\text{B}}$ , and  $K_A$  are the fractional random errors in  $C_{m}$ ,  $C_B$ , and A respectively, and were arbitrarily assigned values of 0.03, 0.03, and 0.06. The standard deviation in the structure factor can then be calculated by the method of finite differences  $(44)$  as

$$
\sigma_{\rm F} = \left(\frac{c_{\rm R}}{\text{L}_{\rm DA}} + \frac{\sigma_{\rm I}}{\text{L}_{\rm P}}\right)^{1/2} - F_{\rm obs}
$$

Where Lp is the Lorentz-polarlzation factor and the observed

structure factor

$$
F_{\rm obs} = \left(\frac{c_{\rm R}}{L_{\rm DA}}\right)^{1/2}.
$$

Consideration of symmetry-extinct data determined that only those reflections with  $F_{\text{obs}} \geq 3\sigma_{\overline{p}}$  would be considered observed and used in the refinement. Of the 1122 measured reflections, 839 met this criterion. In the least-squares refinement the structure factors of the individual reflections were given weights equal to the square of the reciprocal of their respective standard deviations.

## $(B1\frac{5}{9}^+)$ (BiCl<sup>2</sup>-)<sub>2</sub>(Bi<sub>2</sub>Cl<sup>2</sup>-)<sub>1</sub>/<sub>2</sub>

The structure of this compound was solved in **I96I** (17) using film techniques. The least-squares refinement was carried to convergence using only isotropic thermal parameters. Because of the evidence in the final difference map of high anisotropy in the thermal motion of the bismuth atoms, and because of the desire to compare the nonabismuth cations In the two environments, the further refinement of the structure using anisotropic thermal parameters was attempted. Two independent sets of data were used in this process: the original film data and a new set of data collected by counter methods.

 $Bi_{12}Cl_{14}$  was reported to belong to the orthorhombic crystal system with lattice parameters a = **23.057 ± 0.002,**   $b = 15.040 \pm 0.007$ , and  $c = 8.761 \pm 0.003$  Å. Extinction

conditions limited the choice of space groups to Pnnm  $(D_{2n}^{12}$ , No. 58) or Pnn2  $(C_{2n}^{10}$ , No. 34); successful refinement being achieved in the centric group. The intensities of 1957 observed reflections were earlier recorded with Cu  $K_{\alpha}$ radiation using the equi-inclination Weissenberg technique supplemented by precession photographs. The data were treated for Lorentz, polarization, and absorption effects. The further refinement was begun using these reported data  $(45)$ , unit weights, and a different scale factor for each of the six layers, hkO through hk5, of data. 2881 peak height Intensities were measured within one full octant of a two theta sphere of 50° using Mo  $K_{\alpha}$  radiation and an automated Hilger-Watts four-circle diffractometer with a take-off angle of 4.5°. As described elsewhere (4l), the diffractometer aligned the crystal after being given the location of three reflections whose indices collectively included non-zero terms for h, k, and 2. The center of each reflection was determined by left-right, top-bottom beam splitting. As part of the alignment procedure the lattice parameters computed from the orientation matrix were printed by teletype. The values of the unit cell constants determined by this procedure were substantially in agreement with those reported earlier:  $a = 23.067(6)^{\frac{1}{2}}$ ,  $b = 14.992(19)$ , and  $c = 8.772(8)$   $\text{Å}$ .

<sup>^</sup>Estimated errors in parentheses were based on the average deviations among several determinations; hence, they represent precision rather than accuracy.

As described above the new data were converted to integrated Intensities and corrected for Lorentz, polarization, and absorption effects. The transmission factors, calculated for the cylindrical crystal (0.237 x O**.167** mm. in diameter) using a linear absorption coefficient of **678.2,**  ranged from 0.11% to 1.27%.

The standard deviations of the structure factors were estimated as before except the fractional random errors in the total count, background count, and absorption factor were assigned values of 0.10, 0.10, and 0.15, respectively. The square of the reciprocal of the standard deviation was used as the weight of each structure factor in the least-squares refinement. 1705 reflections exceeded the  $3\sigma_{\overline{n}}$  threshold established by consideration of symmetry-extinct data as the criterion for use in refinement.

 $\mathfrak{f}$ 

#### RESULTS AND DISCUSSION

Synthesis and Preliminary Characterization

The isolation of new compounds prepared neat in the molten salt media was hampered by the tendency of the solidified solvent to adhere to the product. Even the determination of the composition of the phases was prevented by the very slow rate at which these systems reached equilibrium. This also made thermal analysis of limited value. Similarly the meticulous work of Corbett (3) failed to establish the true stoichiometry of the bismuth subchlorlde (17) so that the elucidation of the phase relationships in the  $B1-B1C1_3 \cdot 3A1C1_3$ system was only made possible by the multifarious spectral  $(20,21)$ , synthetic  $(19)$ , and crystallographic  $(46)$  investigations. Many of these approaches rely on the availability of a foreign solvent media, such as NaAlCl, and KAlCl,, which allow synthesis and recrystalllzatlon by classical methods and provide a means to gauge the amount of solvent present in the Isolated phase. Unfortunately the alkali metal hexachlorometallate systems appropriate for this investigation do not provide a suitable temperature range to serve as a solvent for the salts of Interest. The attempts to prepare the pure phases are detailed below, but they were largely inconclusive. Resort to crystallography was necessary and possible in the hafnium system, but the lack of suitable crystals prevented this recourse In the other cases. The far-Infrared spectra of

reduced compounds and compositions were gathered to allow a comparison between the synthetic results as well as an evaluation of this technique as a means of identifying the polynuclear cations.

#### The Bi-BiCl<sub>3</sub> .3/2HfCl<sub>4</sub> System

The initial experiments performed with this system were exploratory equilibrations of bismuth metal and a 3:2 mixture of hafnium tetrachloride and bismuth trichloride. The unreacted bismuth was recovered, cleaned, and weighed. The amount of bismuth consumed was taken as the difference between the weights of the metal added and recovered. The results of these experiments are shown in Table 2. Powder patterns were taken of the product of each equilibration. Comparison between patterns from different equilibrations showed them to be identical. The disagreement in the extent of reduction appeared to be due to the difficulty encountered in cleaning the recovered metal pellet. No relationship was evident between the individual results and any of the other experimental conditions. A sample was prepared with 68 mole % bismuth in BiCl<sub>3</sub>.3/2HfCl<sub>4</sub>. The powder pattern of this composition was a superposition of the diffraction patterns of the unknown phase and the starting salts. This suggested that only one reduced compound was formed.

Thermal analysis measurements were made on samples ranging from 50 to 95 mole % bismuth. The thermal halts were





aExtent of reduction is defined as:

 $\mathcal{L}_{\mathcal{A}}$ 

moles Bi (react) x  $100$   $\text{moles }$  Bi (react) + moles  $\text{B1Cl}_3$ 

 $\sim$   $\pm$ 

 $^{\text{b}}$ The weight of bismuth reacted is indicative of the size of the sample.

 $\label{eq:2.1} \frac{1}{2} \int_{\mathbb{R}^3} \frac{1}{\sqrt{2}} \, \frac{1}{\sqrt{2}} \,$ 

 $\frac{1}{4}$  .

generally not reproducible and showed large degrees of supercooling. Although no additional Information was deduced about the phase relationships In this system, a halt at 507° was visually confirmed to be the incongruent melting of the reduced phase. Fortuitously equilibration of excess bismuth and the stoichiometric mixture of  $B1Cl<sub>3</sub>$  and  $HfCl<sub>k</sub>$  just below the peritectic temperature resulted in the growth of acicular crystals from the bulk product. A typical reaction tube with crystals is pictured in Figure 2. Further attempts at determining the stolchlometry of the reduced phase by classical techniques were abandoned. The precise composition of this compound,  $Bi_{10}HF_3Cl_{18}$ , was determined by the X-ray structure determination described below.

The crystal structure revealed the ionic formulation  $(Bi^{\dagger}) (Bi_{\mathfrak{q}}^{5\dagger}) (HfCl_{\mathfrak{q}}^{2\dagger})$ , to be an appropriate description of the primary Interactions in the solid. This raised the possibility of a paramagnetic contribution to the bulk magnetic susceptibility from the  $B1^+$  ion which has a  $6s^2$   $6p^2$  electronic configuration. The magnetic susceptibility measured by the Faraday method eliminated this speculation. Table 3 summarizes the experimental data. The total molar susceptibility,  $X_{M}$ , is -964 ± 34 x 10<sup>-6</sup> emu/mole and is independent of temperature in the range 77° to 298°K. Subtraction of the estimated diamagnetism for the  $B1^{+3}$ ,  $Hf^{+4}$ , and Cl<sup>-</sup> cores (47) leaves a residual molecular diamagnetism of  $-198 \times 10^{-6}$  emu/mole. This high value of the residual dlamagnetism indicates that

 $\frac{1}{3}$ 

Figure 2. Crystals of  $Bi_{10}Hf_3Cl_{18}$  projecting from the bulk sample

 $\sim 10^{-11}$ 

 $\sim 10^{11}$ 

 $\mathcal{L}^{\text{max}}_{\text{max}}$  and  $\mathcal{L}^{\text{max}}_{\text{max}}$






 $a_{\text{Molecular weight}} = 3263.42 \text{ grams/mole.}$ 

^Measurements were made **at** five field strengths and susceptibilities at infinite field were obtained from a leastsquares fit to a Hondo-Owens plot.

 $c$ Values should be multiplied by  $10^{-3}$ .

the twenty-two valence electrons of the  $B1_9^6$ <sup>t</sup> ion make a sizable contribution to the diamagnetic susceptibility. This conclusion is compatible with the highly delocalized valence electron molecular orbital description proposed by Corbett and Bundle (18) for this cation.

The ligand field theory used to Interpret the spectrum of the  $B1^+$  ion in molten salt media (23) predicts a  ${}^{3}P_0$  diamagnetic ground state, in agreement with experiment. This theory can be used to assign the single peak observed at 500 mu (20 kK) in the ultraviolet-visible spectrum of  $Bi_{10}Hf_3Cl_{18}$  to a  ${}^{3}P_{0}$  <->  ${}^{3}P_{2}$  transition. From Table 4 it is apparent that

the AlCl <sub>3</sub> -NaCl eutectic <sup>a</sup>		
$E$ (kK)	$104$ f	
11.1 14.4 15.2 17.1 30.0 32.5	$\frac{0.7}{4.5}$ 5 37 0.3 1.5	

Table  $4$ . Band energies and oscillator strengths for  $B1^+$  in

 $a$ Contents of this table are taken from (20)

the 17.1 kK transition is the most intense in the  $B1^+$  spectrum in the AlCl<sub>3</sub>-NaCl eutectic. Table 5 lists the observed band



Table 5. Band maxima for  ${}^{3}P_{0}$  <->  ${}^{3}P_{0}$ 

maxima for this transition in several molten salts. The energy of this transition seems to shift to higher energy with increasing field strength. Bromide ligands provide weaker fields than chloride ligands as evidenced by their relative positions in the spectrochemical series. The relative basicities of the AlCl,-NaCl and ZnCl,-KCl eutectics implied by this correlation are consistent with the existence of the large low-field  $B1_{8}^{2+}$ ion in the former but not the latter system (21). The assignment made for the spectrum recorded using the Bi-BiCl3 system is questionable. The band was attributed to  $B1^+$  (13) and does have a reasonable position for the band maximum, but the molar absorptivity is high relative to the intensity of this band in other media. Finally the dlnegatlve hexachlorohafnate(IV) ion provides a much stronger field than the mononegatlve anions probably present in the molten salts. The assignment is considered reasonable but not exclusive.

#### The Bi-BiCl, 3TaCl, system

Thermal analysis measurements on this system gave results which were not reproducible except for a halt at  $\sim$ 440°. This temperature was confirmed visually as being the melting point of the most reduced phase. Samples containing 63.1, 80.2,  $83.4$ ,  $90.5$ , and  $93.8$  mole % bismuth in BiCl<sub>3</sub>.3TaCl<sub>5</sub> were prepared by equilibration for 2-1/2 weeks. Earlier thermal analysis had revealed a halt in the range from 310° to 330° for compositions greater than 75 mole  $%$  metal. Since this halt may be due to the melting of an intermediate phase, the first four samples were equilibrated at 300° and the last sample was heated at 410°. The powder patterns of the products are listed in Table 6. The BiCls-TaCls system is a simple eutectic (30), therefore the powder pattern of the unreduced composition will be a superposition of the diffraction lines of the constituent salts.<sup>1</sup> There appear to be two new phases formed in the reduced system. The intermediate phase has a composition between **80** and **83** mole % bismuth and the most reduced phase has a composition between 90 and 94 mole % metal.

 $T$ The principal diffraction lines for BiCl<sub>3</sub> (48) with the relative intensities in parentheses are: 4.75(100), **4.07(18), 3.68(18),** 3.365(33), 2.86(33), 2.73(15), **2.58(24),**   $2.44(21)$ ,  $2.38(17)$ ,  $2.23(32)$ ,  $2.15(14)$ ,  $2.04(22)$ .

The principal lines for TaCls (J. Ebner, **1969,** personal communication) are:  $9.0(50)$ ,  $5.8(100)$ ,  $5.2(80)$ ,  $5.05(80)$ ,  $4.49(90)$ ,  $4.29(80)$ ,  $3.35(25)$ ,  $3.10(30)$ ,  $2.95(25)$ ,  $2.74(100)$ ,  $2.64(75)$ ,  $2.53(75)$ ,  $2.47(75)$ ,  $2.44(30)$ ,  $2.36(30)$ ,  $2.13(100)$ , **2.10(25), 2.03(25).** 

80 80 5 40 10.3 10.3 10.3 10.3 30 30 9.0 10 9.0 9.0 8.7 8.6 8.7 40 20 8.0 30 $\overline{\phantom{0}}$ 7.7 10 7.6 $\blacksquare$ 7.3 10 $\overline{\phantom{a}}$ 5.8 5.8 5.7 20 35 30 5.42 5.4 5.41 40 40 50 $5 - 3$ 5.3 30 5.3 100 100 65 5.2 5.15 40 5.18 50 $4.8 - 4.9$ 4.8 4.80 25 20 30 4.4 4.42 4.4 4.4 - 5 20 20 4.05 4.05 $\begin{array}{c} 5 \\ 60 \end{array}$ 10 $\ddot{ }$ b 3.80 $3 - 80$ 3.81 50 10 b $\overline{\phantom{0}}$ 3.78 70 $3 - 79$ 30 3.7 $5\sigma$ 3.69 $3 - 70$ 60 50 3.66 40 3.60 3.60 20 20 3.50 3.51 5/65 $\overline{5}$ 20 3.50 3.45 3.43 35 35 3.43 $\ddot{\phantom{a}}$ 3.30 50 $3 - 30$ 40 3.30 $3 - 30$ 30 d 30 2.90 85 10 2.92 2.93 10 $\overline{\phantom{a}}$ 2.80 35 $\overline{\phantom{0}}$ 2.77 15 2.73 80 60 2.73 2.72 40 d 2.70 100 2.65 2.64 30 25 2.65 5 2.51 15 2.50 20 2.45 2.46 2.45 50 2.45 40 20 40 2.38 2.37 40 70 $\overline{\phantom{0}}$ $\overline{\phantom{a}}$ 2.33 2.34 45 35 2.33 35 - 2.32 30 $\blacksquare$ 2.28 35 2.28 20 2.20 2,20 15 10 $\overline{\phantom{a}}$ $\overline{\phantom{a}}$ 2.14 60 2.13 100 40 2.14 30 2.15 2.12 2.10 20 20 2.11 10 $\overline{\phantom{0}}$ 2.06 10 10 2.07 10 2.07 5 10 2.02 15 2.03 2.02	30 10.3 $\frac{5}{5}$ 9.0 5 8.7 $\overline{\phantom{0}}$ 5 15 7.3 - 8 5.45 d ▀ 5.22 40 $\frac{1}{5}$ $3 - 77$ 25 3.65 30 8 3.60 $\overline{\phantom{0}}$ $\overline{a}$ 3.42 75 $3 - 30$ 30 2.92 100 2.80 5 $\rightarrow$ 100 2.70 2.64 15 40 2.50 2.45 60 $\rightarrow$ 2.33 75 2.31 30 $\rightarrow$ 2.41 80 2.08 10 2.01 20

Table 6, Powder patterns of reduced compositions in the Bi-BiCl<sub>3</sub> · TaCl<sub>5</sub> system

a The composition given in mole % bismuth in  $B1Cl<sub>3</sub> \cdot 3/2HfCl<sub>4</sub>$ .

 $b_{b=broad}$ , d=diffuse.

 $\bar{z}$ 

Microscopic examination of one of the thermal analysis samples had shown the presence of crystals with two different colors, dark red and black (D. A. Lokken, 1969, personal communication). The red color appeared similar to the intermediate phase found in the Bi-BiCl<sub>3</sub> . 3AlCl<sub>3</sub> system. The two reduced compounds in the aluminum chloride pseudobinary system have compositions within the ranges determined for those compounds formed in the tantalum chloride system;  $B15(A1C14)3$ with 80 mole % bismuth and Bi<sub>k</sub>AlCl<sub>k</sub> with 91.7 mole % metal. Although the similarity is noted between the systems containing tantalum pentachlorlde and aluminum chloride, there is no unambiguous evidence to support the relationship.

Because difficulty was encountered in attempting to prepare the pure phases by direct synthesis, the most promising procedure for Isolating the phases was considered to be recrystallization from a foreign solvent. Since the alkali metal chloride-tantalum pentachlorlde systems do not have suitable melting points, an attempt was made to use the ammonium analogue which is molten in the temperature range of interest. Unfortunately the work of Morozov and Toptygin (49) may be correct and the fused solvent was Intensely colored. The other possibility is that the ammonium ion reduced the TaCl<sub>6</sub> to TaCl<sub>6</sub><sup>-</sup> which is highly colored (50). This hindered the observation and Isolation of the desired compounds. Further work in this system was discontinued. A source of the Bi<sub>3</sub><sup>+</sup> ion in the solid state was sought. It was hoped that the peculiarities of the

Interactions of this cation and anions In the solid might give some Insight as to why Its existence seems limited to media containing mononegative anions. The compound  $Bi_1 \rho Hf_3Cl_{18}$  has the same average oxidation state for bismuth as the  $B15^{3+}$ cation but the species present in the solid are  $B1^+$  and  $B1_9^{5+}$ ions.

## The B1-B1Br<sub>3</sub> AlBr<sub>3</sub> system

The compound  $B1_5(A1Cl_4)$ , has been shown (46) to crystallize in a twinned pseudo-orthorhombic cell. The investigation of the Bi-BiBr<sub>3</sub> AlBr<sub>3</sub> system was conducted in the hope of obtaining suitable single crystals for an X-ray crystallographic study. The possibility existed that the change in anion might result in a different crystal habit. No good crystals were found. The powder patterns of the two compounds, shown in Table 7, suggest that the two phases may be isostructural. The tetrabromoalumlnate obviously must have a larger unit cell. The relative intensities are dependent on the diffracting power of the constituent atoms; therefore the difference in relative intensities between the chloride and bromide salts are reasonable.

### The B1-B1Cl<sub>3</sub>-KCl system

The chemical composition of the bismuth subchlorlde has not been determined accurately by chemical analysis. The empirical formula Bi6Cl7 was determined by the elucidation of

$B1_5(A1CL_4)_3^a$			$\text{Bi}_5(\text{AlBr}_4)$ <sub>3</sub>	
$d(\lambda)$	$\underline{\text{I/I}}_{\underline{\textbf{0}}}$	$\underline{d(\overset{\mathtt{o}}{\mathtt{A}})}$	$I/I_0$	
9.45 8.60 7.45 6.8 6.58 5.97 4.96 4.84 3.83 3.29 3.25 3.103 3.051 3.025 2.966 2.914 2.692 2.657 2.569 2.547 2.369 2.345 2.302 2,266 2,241 2.210 2.189 2,164 2,102 1.998 1.833 1.710 1.670 1.647 1.500 1.465 1.413	0.1 100 10 10 0.1 70 30 50 50 10 30 0.5 10 10 40 10 50 50 30 60 0.5 0.5 10 10 10 0.1 0.5 15 diffuse 15 diffuse 10 0.5 50 10 20 10 30 10	9.8 8.8 7.5 6.25 5.1 3.95 3.57 3.25 3.12 $2.9 - 3.0$ 2.78 2,66 2,44 2.14 2.08 2.03 1.88 1.85 1.80 1.76 1.73 1.70 1.64 1.535 1.51	3 100 5 50 85 8 8 5 20 100 35 30 8 8 20 15 10 10 15 5 10 15 15 10 5	

Table 7. Interplanar spacings and relative intensities for diffraction lines of  $Bi_5(A1Cl_4)$ , and  $Bi_5(A1Br_4)$ ,

Diffraction Unes for this phase are reported by Corbett (19) and Levy, et al. (16). The pattern reported here was taken from Qrossalnt (51) who placed the relative intensities on one scale.

the crystal structure. In an attempt to substantiate the results of the X-ray study, samples of the lower chloride were prepared from Bi-BiCl<sub>3</sub>-KCl melts and the potassium content of the crystals used to correct their analysis for solvent retention. It Is Important to realize that correction for solvent retention can only be valid If the solvent retained by the crystals has the same composition as the solvent analyzed after having removed the crystals. This condition can only be approximated if the gross composition of the solvent does not change appreciably during crystallization. The correction will also be more accurate the greater the difference in composition of the solvent and the crystals. Table 8 shows the results of a typical analysis. The magnitude of the "solvent correction" suggests that the procedure is not capable of the necessary accuracy to differentiate between BiCl (85.85 weight % bismuth) and  $BI_{12}Cl_{14}$  (83.47 weight % bismuth).

Although the attempted analysis was not successful, experiments did reveal an improved synthetic procedure for the reduced compound. Contrary to the usual arguments about acidbase interactions, the addition of the base potassium chloride to the bismuth-bismuth trichloride binary increased the quantity of crystalline  $Bi_{12}Cl_{14}$  recovered from the melt. Further experiments revealed that whereas small additions of the base to the binary increased the reduction, large concentrations of KCl prevented any reduction from occurring. These observations can be rationalized in terms of the structure of

**!|2** 



Table 8. Typical analysis of  $Bi_{12}Cl_{14}$  grown from Bi-BiCl<sub>3</sub>-KCl melts.

 $a$ Computed using the formula employed by Corbett (19).

<sup>b</sup>It is important to note that disproportionation and subsequent loss of CI may have occurred in an attempt to remove excess BiCla before analysis, giving high values for Bi content.

BiCl wt % Bi=85.85%

 $B1_{12}Cl_{14}$  wt %  $B1=83.47%$ 

average wt  $\cancel{\epsilon}$  B1(Corbett) =  $84.93\%$ <sup>b</sup>

the KCl-BlCla melts. Phase (52), cryoscoplc (53), and spectral (54) studies suggest the presence of BiCl<sub>5</sub> and BiCl<sub>5</sub><sup>-</sup> ions in the melt with very little BiCl<sup>3</sup>. Since these ions are present in the solid  $BI_{12}Cl_{14}$ , it appears that a common ion effect Increased the amount of the phase of interest. At higher chloride concentrations most of the bismuth in the melt is complexed thus increasing the stability of the  $Bi^{3+}$  state and decreasing the reduction to the  $B15^+$  ion. The best preparative conditions appeared to be reduction with excess bismuth of melts containing 20 to 30 mole % KCl in BiCl<sub>3</sub>.

# The far-infrared spectra

The far-infrared spectra were recorded for  $BI_{10}Hf_3Cl_{18}$ ,  $B1_{12}C1_{14}$ , and several reduced compositions in the Bi-BiCl<sub>3</sub>•3TaCl<sub>5</sub> system. Traces of the spectra are pictured in Figures 3, 4, and 5, respectively, and in Figure 6 for the compounds  $Bi_S(AICl_*)$ <sub>3</sub> and  $Bi_*(AICl_*)$ . Possible bands have been marked with arrows. Representative bands for the stoichiometric compounds are listed in Table 9. Table 10 shows the bands for the reduced compositions. Table 11 contains the known vibrational frequencies of the anions present in the reported phases. The assignment of bands arising from anion vibrations are discussed below, but

 $1$ The spectra of the tetrachloroaluminate salts were generously furnished by Dr. D. J. Prince (D. J. Prince, 1969, personal communication).

Figure 3. Far-infrared spectra of  $Bi_{10}Hf_3Cl_{18}$ 

 $\sim 10^{-1}$ 

 $\bar{z}$ 

 $\Delta\sim 1$ 

 $\mathcal{L}^{\text{max}}_{\text{max}}$  and  $\mathcal{L}^{\text{max}}_{\text{max}}$ 



 $\overline{a}$ 

l,

 $\overline{\phantom{a}}$ 

 $\bar{\bar{z}}$ 

 $\bar{\boldsymbol{s}}$ 

Figure 4. Far-infrared spectra of  $B1_{12}Cl_{14}$ 

the contract of the contract of the

 $\label{eq:2.1} \mathcal{L}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}})) = \mathcal{L}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}})) = \mathcal{L}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}))$ 



Figure 5. Par-infrared spectra of reduced compositions in the system  $B1-B1C1_3 \cdot 3TaC1_5$ :

- (a) **80.2** mole % bismuth;
- (b) a different sample with the same composition as (a);
- (c) sample (b) rerun 4 hr. later;
- (d) 90.5 mole % bismuth; and
- $(e)$  93.8 mole % bismuth



Figure 6. Far-infrared spectra of (a)  $B1_5(A1Cl_4)$ , and  $(b)$  Bis $(AICI<sub>4</sub>)<sub>2</sub>$ (D. J. Prince, **1969,** personal communication)

 $\sim$ 

 $\sim$ 

 $\sim$ 

 $\mathcal{L}$ 



 $\sim$ 

Far-infrared frequencies of Bi10Hf3Cl18, Bi<sub>12</sub>Cl<sub>14</sub>, Bi<sub>5</sub>(AlCl<sub>4</sub>)<sub>3</sub>, and Bi<sub>4</sub>(AlCl<sub>4</sub>) in the solid state (in cm<sup>-1</sup>) Table 9.  $\sim$   $\sim$   $\sim$ 



 $a_{s=strong,$  m=medium, w=weak, v=very, sh=shoulder, br=broad, asym=asymmetric.

bAssigned bands are for the anion

 $\mathcal{L}$ 



 $\label{eq:1} \mathcal{L} = \mathcal{L} \times \mathcal{L}$ 



a<sub>The composition is given in mole % bismuth in</sub>  $B1C1_3$   $\cdot$  3TaCl  $_5$  .

 $^{\text{b}}$ s=strong, m=medium, w=weak, v=very, br=broad, sh=shoulder.

consideration of the metal-metal vibrations is complicated by the poor spectra and the possible presence of lattice modes in the region below 100  $cm^{-1}$ . The region anticipated for bismuthbismuth vibrations is between 50 and 200  $cm^{-1}$  based on the known frequency of the gaseous dimer (55) and the estimated frequencies of the gaseous trimer (linear) and tetramer (tetrahedral) (56).

The octahedral anions have six fundamental bands;  $v_1(a_{1g})$ ,  $v_2(e_g)$ , and  $v_5(t_{2g})$ , are expected to be Raman active, while  $v_3(t_{1u})$  and  $v_4(t_{1u})$  are infrared active and  $v_6$ is vibrationally inactive. In both the HfCl<sup>2</sup> $\overline{c}$  and TaCl<sup>7</sup><sub>6</sub> systems the vibrationally active fundamental modes appear in the infrared spectra. The assignment of almost all bands at frequencies greater than  $13^{1}$  cm<sup>-1</sup> to vibration of the anions seems reasonable on the basis of intensity and coincidences with the known frequencies. The appearance of the  $v_1$ ,  $v_2$ , and V5 modes in the infrared indicate that the ions are distorted from  $O_h$  symmetry. The intensities suggest strong interactions between the cations and anions in the lattice. These effects are also noted with the tetrachloroaluminate salts where the triply degenerate  $v_3$  and  $v_4$  modes are each split into two bands.

The very low intensity of the other bands indicate that the infrared spectra are not very useful in fingerprinting the different phases. It is interesting to note change in spectra for the different compositions in the  $Bi-BiCl_3 \cdot 3TaCl_5$  system.

Octahedral ions						
$HfCl62$ = :						
Compound	$v_1(r)$	$v_2(r)$	$v_3(1)$	$v_{4}(1)$	$v_{5}$ $(r)$ $\frac{V}{4}$	Reference
$(Et4N)2 (HfCl6)$	333	237	288	145	$(80)^a$ 157	(57)
Aqueous solution saturated with HCl	331 $s^b$	$204$ m	$272 \text{ s}$ 304 <sup>}</sup> sh	$\overline{\phantom{m}}$	152 m	(58)
$(Et2NH2)2 (HfCl6)$	$328$ vs	264 w	$286$ <sub>1</sub> s $147$ <sub>1</sub> s $273$ <sup>1</sup> s,sh $138$ <sup>1</sup> s,sh		$\frac{153}{163}$ }ms	(59)
$(Et_hN), (HfCl_h)$	$326$ s	$257$ w.sh	$275$ s	145 s	156 s	(60)
Cs <sub>2</sub> HfCl <sub>6</sub>	333 s	(261)	284 s	150 s	$167$ s $(110)$	(60)
	333 mw	235 m	280 s	$134$ ms	$154$ sh <sub>}</sub> 165 ms <sup>}</sup>	this work

Table 11. Fundamental vibrations of the anions present in polynuclear bismuth compounds  $(in cm^{-1})$ 

 $a$ <sub>The</sub> frequencies listed in parentheses were not observed but were calculated from combination bands.

 $b_s$  = strong, m = medium, w = weak, v = very, sh = shoulder.

 $\sim 100$ 

 $\sim 10^{11}$  km  $^{-1}$ 



 $\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}})$  and  $\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}})$  and  $\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}})$ 

The results appear to corroborate the tentative conclusions reached earlier on the basis of the powder pattern data; there may be two reduced compounds in the system. Unfortunately the poor quality of the spectra preclude more definitive conclusions.

The Crystal Structure of  $(B1<sup>+</sup>)(B1<sup>5+</sup><sub>9</sub>) (HfCl<sup>2-</sup><sub>6</sub>)<sub>3</sub>$ 

## Solution and refinement

The structure was solved by conventional heavy atom techniques (34). An unsharpened Patterson map with approxi-0 mately 0,25 A resolution and sectioned along the W-axis was computed using the observed data. A notable feature of the Patterson function was the periodicity of vector maxima lying in UV planes for which values of W occurred in multiples of one-sixth. Maxima along the W-axis were found at (0,0,1/3) and (0,0,2/3) in addition to the origin peak. Analysis of the Patterson proceeded using Marker lines and planes calculated from the general positions of the centric space group. Centrosymmetry had been suggested by the Howells, Phillips, and Rogers statistical test (62).

Using the Barker peaks and the inter-set vectors two compatible sets of bismuth atoms, one twelve-fold degenerate and one six-fold degenerate, were found which described a Big

unit. These positions were refined by least-squares<sup>1</sup> minimization of the function  $R_{\omega} = [\sum \omega (|\vec{F}_0| - |\vec{F}_C|)^2/\sum \omega |\vec{F}_0|^2]^{1/2}$  to yield a conventional residual  $(R_1 = \frac{\Gamma(|F_0| - |F_0|)}{F_0})/|F_0|$  of 0.32, The structure factors phased by this model were used to construct a Fourier synthesis which revealed the hafnium  $(R_1 = 0.30)$  and subsequently the chlorine  $(R_1 = 0.20)$  positions In a HfCle grouping. A Fourier difference map computed with the resultant model with the stoichiometry  $Bi_{1B}HF_{6}Cl_{1B}$  showed electron density along the Z axis at  $z \approx 0.1$  and at the symmetry related positions.

The symmetry of the centric space group requires atoms at  $0,0,1/2-z$ ,  $0,0,1/2+z$ , and  $0,0,-z$  when an atom is present at  $0,0,2$ . Since the Z axis is only 10.692  $\lambda$ , to prevent two atoms from being in too close proximity to each other either this site must be partially occupied or the acentric space group, with symmetry generating only positions  $0,0, z$  and  $0,0,1/2+z$ , must be correct.

Because of the nature of the least-squares equations one cannot take a centric model and refine it In an acentric space group (65,66). The results obtained are unreliable because of the correlation of parameters related by the pseudosymmetry.

 $<sup>1</sup>A$  full-matrix least-squares refinement was carried out</sup> using ORFLS (38) modified to correct for both real and imaginary parts of anamolous dispersion. The scattering factors used were taken from the tables of Hansen, et al.  $(63)$ and corrections for anomalous dispersion were taken from the International Tables (64).

An attempt, albeit questionable, was made to avoid this problem by fixing the parameters of the atoms In the model, and to refine only the Bi atom added to. the structure in the alternative ways; i.e., in one of the two independent positions in the acentric space group or with partial occupancy of the position in the centric space group. The conventional R factor for refinement in the centric model, 0.163, was lower than that computed for refinement in the acentric model, **0.171.** Further refinement of the trial structure in the centric space group with 50% statistical occupancy of the axial bismuth position converged to  $R_1 = 0.157$  and  $R_n = 0.167$ .

At this stage three reflections, (100), (010), and (001), for which  $\theta < 2.25^{\circ}$ , were found to be unreliable because in each case the diffracted beam had passed through the metal casing of the scintillation counter. This gave rise to a nonlinear background count and required discarding these data. The remaining data were used for refinement in which the temperature factors of all the atoms were allowed to vary anisotropically  $(R_1 = 0.102; R_0 = 0.116)$ . These data were then examined for reflections which may have been measured incorrectly. Thirteen additional pieces of data were discarded because they met the criterion  $|F_0 - F_C| \geq 9\sigma_F$ . Further Inspection of these data revealed a common factor; the transmission factors for all these data were less than **1%,**  Successive refinement with the 823 remaining reflections converged to a conventional residual of 0.093 and a weighted

R factor of **0.096.** The standard deviation of unit weight  $[(\sum_{i=1}^N |\mathbf{F}_i| - |\mathbf{F}_i|]^2 / (NO-NV)^{1/2})$  where NO is the number of observations (823) and NV is the number of variables (55)] was **1,87** electrons. This Indicated that the values of the constants chosen to describe the random errors which have subsequently been used in calculating the standard deviations of the individual reflections were too low. The weights were therefore changed by plotting  $\omega \Delta^2$  versus  $F_0$  for 15 overlapping groups of 110 reflections per group and adjusting the weights such that  $\omega \Delta^2$  for the groups were constant  $\int \omega = 1/\sigma^2$  and  $\Delta^2 = (|F_0 - F_c|)^2$ . Refinement using this weighting scheme did not change the conventional R factor, in fact  $R_{\omega}$  increased from **0.096** to **0.107,** but it did lower the standard deviations of the computed distances and angles. In the final least-squares cycle each variable shifted an average of 0.02 times the error associated with It. The largest residuals in the final difference map were  $\pm$  4 e/ $A^3$  in the region of the Bi<sub>3</sub><sup>+</sup> ion and  $\pm$  1 e/ $\AA$ <sup>3</sup> in the region of the B1<sup>+</sup> ion on a scale of 170 e/ $\AA$ <sup>3</sup> for a bismuth atom. The final standard deviation of an observation of unit weight was O**.98** electrons. Based on the agreement of the large structure factors, no extinction correction was necessary.

The final positional and thermal parameters are listed in Table 12, The mean thermal amplitudes of displacement along the three principal axes were listed In Table 13 for the eight



**Table 12. Final positional and anisotropic thermal parameters®" (errors in least significant figures in parentheses)** 

 $\text{exp}[-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}k^2 + 2\beta_{12}hk + 2\beta_{13}hk + 2\beta_{23}kh)].$ 

**^Fractionally (0,5) occupied.** 

Table 12 (Continued)

 $\mathcal{L}^{\pm}$ 



 $\sim$ 

 $\sim 100$  km  $^{-1}$ 

 ${}^{c}$  $\beta$ 's are x 10<sup>4</sup>.

$$
d_{\beta_{22}} = \beta_{11} \text{ by symmetry.}
$$
  

$$
e_{\beta_{12}} = 1/2 \beta_{11} \text{ by symmetry.}
$$





63

 $\bullet$ 

 $\mathcal{L}^{\text{max}}_{\text{max}}$  , where  $\mathcal{L}^{\text{max}}_{\text{max}}$ 

 $\bar{\mathcal{A}}$ 

 $\mathcal{A}^{\mathcal{A}}$ 

 $\sim$ 

 $\label{eq:2.1} \frac{1}{\sqrt{2}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2.$ 

atoms in the asymmetric unit. The observed and calculated structure factors are shown in Figure 7.

## Description and discussion

Ihe crystal structure found for bismuth(I) nonabismuth hexachlorohafnate(IV) is shown in projection in Figure 8 and in a stereogram in Figure 9. The structure is comprised of the large  $B1_3^{5+}$  cations and HfCl<sub>6</sub>" anions which appear to govern the packing arrangement and the small  $B1^+$  ions which restore charge neutrality to the compound and contribute to the lattice energy, and thus to the stability, of the phase. The bonding distances and angles are listed in Tables 14 and 15. The important interatomic non-bonding Interactions are shown in Table 16.

The  $B1a^{5}$  ion is a tricapped trigonal prism similar to that found in bismuth subchloride (see Figure 10). The polyatomic cation has crystallographic  $c_{\text{sh}}$  symmetry. There are four bismuth-bismuth bonding distances in the cation: 3.241(3)  $\AA$  within the triangular face of the prism, 3.737(4)  $\AA$ for the height of the prism, and  $3.086(3)$  and  $3.103(3)$  Å from the prism to the waist atoms. The difference between the two prism to waist distances is Just significant at the level of four times the standard deviation and it is only by this difference that the cation deviates from full  $D_{ah}$  symmetry.

6k

Figure 7. Observed, and calculated structure factors for  $Bi_{10}Hf_3Cl_{18}$  (\* mark reflections not used in final refinement)

 $\mathcal{L}^{\text{max}}_{\text{max}}$  and  $\mathcal{L}^{\text{max}}_{\text{max}}$ 

 $\label{eq:2.1} \frac{1}{\|x\|} \leq \$ 

2222 - 2222 - 2322 - 2322 - 2322 - 2322 - 2322 - 2322 - 2322 - 2322 - 2322 - 2322 - 2322 - 2322 - 2322 - 2322 -<br>2322 - 2322 - 2322 - 2322 - 2322 - 2322 - 2322 - 2322 - 2322 - 2323 - 2323 - 2333 - 2343 - 2343 - 2343 - 2343 Orraaside erberuldistel alubade maarin seistel lindragade ostel produkte sei ,我们的人们的,我们的人们的学习,我们的学习,我们的学生,我们的学生,我们的学生,我们的学生,我们的学生,我们的学生,我们的学生,我们的学生,我们的学生,我们的学生,我们的学生,我们的学生,我们的学生 医神经病 医神经病 医心包膜炎 医心包膜炎 医心包膜炎 医血管下动脉 医心包膜炎 ini nining a a a a aghunda hiding diga basang mang masa sa sa tagatan tanggan tida agusan turunnan minit a a b<br>Tagan matan a disama da intang diang tahun batan missassi sa dina magtama di mathin at dit awaw mbigan dinatan ાં કો અન્યાં કરે છે કો કો કોઈએ. એ . આ એ સ્થાની પ્રાપ્ત માટે કરા કરે . આ તમની દૂધની તેઓના સમાનવ સિનાનો કોઈએ છે.<br>૧.૧૬ પછી કરી વાર જોવાની સમાનવ પર્વત વાર જોન્ની સમાનવ સમાનવ સાથે મારીના માટીની મારીના સમાનવ સાથે સમાનવ વાર માટે ្រាប់ ស្រុក ស្<br>ស្រុក ស្រុក ស្ as State See was as State Valle Lands Back of the and the second state of 2000 and 2000 and 2000 and る はこともちゃりゅしゅっぽー ちゅうしょうきゅうりょうきゅうきょうきょうしょうきゅうじょうしょ しょうしょう しゅっしょう ちゅうじゅつ はいはしこうしょう ウリスト・ライク・コーヒート しゅうしょうしょう しゅうしゅうしゅ ちゅうしゅ ちゅうしゅ じゅん じゅんじん しょうじゅつじょう しょうしょうしょう しょうじょう ところは、そのようなのかないと思います。 このことをしていることはない。 さずなんない とっかい こうしょう しょうどうじょう こうじょう こうしょう いんしょう こうこう こうしょう いっと しょうこう こうしょう こうこう しゃしん しゅうしょう しゃしょう しょうしょう しょうこうしゅん しょうこうしょう יווי את הנו ויי לא היה את היה היה לא היה<br>היה לא היה לא היה<br>ה ואקלה מולה שבון בלאמד לאחר להואה ונמאב ועול.<br>ממורי לספיטמים **להפשם אינו להיות ל**מורי לא להיות הוא להיות להיות להיות להיות היותר המאורים להיות להיותר מה להו<br>מוריי לא ספיט לפיט לשום האל לשל מה הממורי שהים טומית החל לש ללא א - ますし、時、そのできのでまできょうか?自分 注記されるみのでのまでまできるちゅう 自分 問注するでんち 立て ロゴミさんち カアらの まえをんろう しんかんしょうきょう とします。、ここまで、どち、しょう。コ<br>ウェルボッチャウメルドメーションです。これから、これは、これは、これは、これは、これは、これは、これは、これは、これは、これについてのようになります。これは、これはどうなどのようなどのようなのでも、これは、これは、これは<br>いっとう っこりない ちょうてんきょう そうのうか じゃくてん こうきょうかん うちょうせい しゅうしょうしん しょうしん こうていきょうこう ואות האופן האו<br>המוסיקה האופן האופן<br> - ままたてもれるときます。それは、それは、これはこのようには、それはあることについている。 しょうしゅう こうほうせい せいしょうしょう こうじょう こうのはいます こうしょうかん こうしょう アイスター こうしょう こうしょう こうしょう こうしょう しょうしょう こうしょう こうしょう こうしょう こうしょう こうしょう e a contra la construcción de la constitución de la constitución de la constitución de la constitución de la c<br>La constitución de la constitución 

Figure 8. Projection of four adjacent unit cells between  $z=0$  to  $z=1/2$  of  $B1_{10}Hf_3Cl_{16}$  along the (001) direction (single atoms represented by one circle are at  $z=1/4$ , double atoms represented by two concentric circles are related by the mirror plane at z=l/4)


Figure 9. Stereoscopic view of the contents of one unit cell of  $BI_{10}HF_3Cl_{16}$ 

 $\sim 10^{11}$ 



 $\ddot{\phantom{a}}$ 

 $\mathcal{L}^{\pm}$ 



 $\hat{\mathcal{A}}$ 

Table 14. Bond distances and angles<sup>a</sup>

 $\bar{z}$ 



^Standard deviations are given in parentheses and correspond to the least-significant figure. They were calculated by ORFEE (39) using the full variance-covariance matrix produced by the least-squares refinement.

 $\sim$ 

 $\sim 10^{-1}$  km

and a series

 $\sim$ 



Table 15. Distances within ionic entitles calculated with final positional and thermal parameters using the riding model<sup>a</sup> (in  $\chi$ )

a<sub>Interatomic</sub> distances averaged over thermal motion. Second atom is assumed to ride on the first. The function is

$$
\overline{R} = R_0 + (\overline{r_2^2} - \overline{\xi_2^2} - \overline{r_1^2} + \overline{\xi_1^2})/2R_0
$$

where  $R_0$  is the uncorrected interatomic distance,  $r_1^2$  is the mean square radial thermal displacement of atom  $i$ , and  $\xi$  is the mean square component of displacement of atom i in the direction defined by the Interatomic vector.



Table 16. Non-bonding interactions<sup>a</sup> (in  $\lambda$ )

.<br>The contract of the contract of

Angle  $Cl(4)-B1(3)-Cl(4) - 119.7^{\circ}$ 

<sup>a</sup>Standard deviations are given in parentheses and correspond to the least-significant figure,

bInterionic interaction.

Intraionic interaction.

Figure 10. The  $\overline{Bi_g}^+$  cation found in the structure of  $\overline{Bi_1}_0 \overline{Hf}_3 \overline{Ci_{16}}$ 

 $\overline{a}$ 

 $\sim$ 



 $\epsilon$ 

The high symmetry observed for the cation in the solid reflects the absence of any close chlorine-bismuth contacts, the shortest being  $3.374(9)$  Å between Bi(1) and Cl(3).

The cluster geometry is analogous to the ligand framework found In many discrete nine-coordinate metal complexes. The obvious difference lies In the absence of a central atom binding the cluster and In the substitution of bonding interactions for the non-bonding llgand interactions present in the complex. Additionally the prism of the cation is shortened by **6%** from the idealized arrangement of nine atoms equidistant from a central atom. This is the direction of distortion predicted by Kepert **(67)** for coordination compounds from consideration of llgand-llgand repulsions. The explanation of the relative distances in the polyhedron must consider bonding forces. The molecular orbital diagram and the symmetry-adapted basis orbltals used by Corbett and Bundle **(I8)** to rationalize the electronic structure of this cation are pictured in Figure 11. The height of the prism relative to the other distances in the polyhedron must be dictated by the maximization of the waist to terminal orbital overlap. The population analysis of the computed molecular orbltals has shown these interactions to contribute substantially to the overall bonding.

 $\mathbf{v}$  .

The  $B1^*$  ions form a broken column along the Z axis (see Figure 12). Three chlorine nearest neighbors form a plane normal to the Z axis with its center displaced 0.18 Å from the

Figure 11. Molecular orbitals for  $Bi\frac{5}{9}$ <sup>+</sup>

- (a) Symmetry orbitals for Bi<sup>5+</sup> and
- (b) Energy level diagram for  $Bi\frac{5}{9}$ <sup>+</sup> (taken from (l8))

 $\sim 10^7$ 





Figure 12. The  $BI<sup>T</sup>$  ions and their environment in BiioHf<sub>3</sub>Cl<sub>is</sub> (Dotted lines between shaded chlorine atoms delineate the triangular arrangements described in the text.)

 $\ddot{\phantom{a}}$ 

 $\sim 1$ 

 $\mathbb{R}^2$ 



BI + AND ITS ENVIRONMENT IN (BI) (BI9) (HFCL6)3

Bi<sup>+</sup> ion (CL4A about BIl in Figure 12). The distance from the bismuth ion to these chlorine atoms is  $3.188(12)$   $\beta$ , close to the sum of the estimated van der Waals radii, 3.2  $\hat{A}$  (17). There are six next nearest neighbors in triangular units above and below the bismuth ion (CL4B at 3.786(16) and CL2 at  $3.593(19)$   $\AA$  from BIl in Figure 12). The latter pair of triangles are nearly equidistant from the unit formed by CL4A. This results from the accidental equality of the x and y coordinates of CL2 and CL4B and the nearly equal spacing between successive triangles along the Z-axls: CL2 (0,260, 0.070,0.250), **CL4A (0.083,0.824,0.088),** CLUB (0.259,0.083, 0,912), These triangles of chlorine atoms leave an open column along the Z axis, therefore, uniaxial ionic conduction parallel to this axis might be expected.

The symmetry about the bismuth(I) ion is exactly  $C_3$  with only small deviations from  $D_{3h}$  as described; the coordination is best described as substantially trigonal since the CL4A-BI1-CL4A angles are 119,7°. The large anisotropic thermal motion of the bismuth ion can easily be understood in terms of the local environment about the ion. There is another factor which might contribute to the large root-mean-square thermal displacement along the Z axis for this ion,  $0.541(18)$  A. The statistical occupancy of the position is all that is sampled by the X-rays; there exists the possibility that the refined position represents a space-averaged rather than a timeaveraged distribution. This would mean that the individual

bismuth ions do not undergo the motion implied by the thermal parameter. The environment about the Bi<sup>+</sup> ion is clearly pictured in the stereogram in Figure 13.

The hexachlorohafnate(IV) anions are distorted octahedra in agreement with the interpretation of the far-infrared spectra presented earlier. The hafnium-chlorine bond distances are slightly longer than the 2.33  $<sup>0</sup>$  found by electron diffrac-</sup> tion as the mean metal-halide distance in the gaseous  $HfCl<sub>4</sub>$ (68). Surprisingly the two chlorine atoms nearest the relatively high field  $B1^+$  ion have the shortest, and identical, hafnium-chlorine bond lengths, 2.406 %, The longest bond distance in the HfCl<sup>2</sup><sup>-</sup> ion is between the metal and the chloride which is in closest proximity to the large  $B15^+$  ion. This cation is electrophilic, but the distance and the absence of any appreciable distortion of the cation suggest that there are no very strong forces on the ion. It is possible that the high field  $B1^+$  ion polarizes the anion and accumulates charge density in the region between the hafnium and the neighboring chlorides but for the moment the variations in the bond distances within the hexachlorohafnate(IV) ion remain attributed to packing.

The Refinement of  $(B1_9^{5+})\left(B1C1_5^{2-}\right)_2\left(B1_2Cl_8^{2-}\right)_{1/2}$ 

Concomitant refinements were'carried out on the structure of  $Bi_{12}Cl_{14}$  using either the older film data or the new counter data. If a complete anisotropic refinement is attempted with

## Figure 13. Stereoscopic view of the environment about the  $B1^+$  ions in  $B1_{10}Hf_3Cl_{18}$

 $\mathbb{R}^2$ 

 $\sim 10^7$ 

 $\sim 10^{-1}$ 





film data collected about only one axis and a different variable scale factor is allowed for each layer, the leastsquares equations will generate a singular matrix. In order to determine a common scale factor for the layers hkO to hk5, several cycles of least-squares were therefore performed using isotropic thermal parameters and unit weighting of each observed reflection. After the data had thus been converted to a single scale, the thermal parameters of the bismuth atoms were allowed to vary anisotroplcally. The weights of the reflections were subsequently adjusted using the program OMEGA (C. Hubbard, 1969, personal communication). Successive Iterations of least-squares yielded a conventional residual of 0.148 and a weighted residual of 0.206, Attempts to vary the chlorine atoms anisotroplcally gave non-positive definite, i.e., physically unreal thermal parameters. Since the earlier Isotropic refinement using these data with seven-tenths as many parameters had resulted in values of O**.15I** and O**.I86** for  $R_1$  and  $R_n$  respectively, it appears that the adjusted weighting scheme used was not as good as the weights chosen by the previous workers to reflect the experimental uncertainties in the data. A difference Fourier synthesis revealed peaks of  $\pm$  9 e/ $A^3$  in the neighborhood of the B1<sup>5</sup> cation and  $\pm$  4 e/ $A^3$ in the region of the chlorine atoms on a scale of **I86** and  $35 e/\overset{\circ}{\text{A}}$  for the bismuth and chlorine atoms respectively.

Least-squares refinements using counter data yielded values of  $R_1 = 0.208$  and  $R_4 = 0.232$  with isotropic thermal

parameters and  $R_1 = 0.169$  and  $R_n = 0.192$  with anistropic thermal parameters for the bismuth atoms. It was found that the chlorine thermal parameters could not be varied anisotropically. The poor quality of these data were manifest in the observation that twenty-one symmetry-extinct data had structure factors greater than three times their Individual standard deviations. To confirm that this did not indicate that the extinction conditions were in error, the symmetry-extinct data were carefully examined. It was found that the net intensity of these data and the subsets of these data represented by the hOl and Okl zones were arranged in a Gaussian distribution about zero counts.. This suggested that the troublesome peaks were only part of large random fluctuations. The final difference map revealed peaks of  $\pm$  5 e/ $\lambda^3$  near the bismuth positions and  $\pm 6$  e/ $\AA$ <sup>3</sup> near the chlorine positions on a scale of 195 and 28  $e/A^3$  for bismuth and chlorine respectively. The relatively large peaks near the chlorine atoms show the sensitivity of the lighter atoms to the quality of the data and an evidently poor absorption correction. Structure factors from symmetry-equivalent data differed from each other by an average of **1S%.** 

The structure of  $BI_{12}Cl_{14}$  is shown in projection onto the (001) plane in Figure l4. The atom positions from the three different refinements are shown in Table 17. The bond distances within the  $B14^+$  cation computed from these parameters are shown in Table IB,

Figure 14. Projection along the (001) axis of one octant of the structure of  $B1_{12}Cl_{14}$  (taken from (17))

 $\sim$ 

 $\sim$ 

 $\sim 10^{-1}$ 

 $\sim$ 







^Initial refinement of film data isotropic thermal **parameters** (71 variables, **1957** data, **27.6** observations/parameters) [from **(17)].** 

^Refinement of film data using anisotropic thermal **parameters** for bismuth atoms (99 parameters, 1957 data, 19.9 observations/parameters).

^Refinement of counter data using anisotropic thermal parameters for bismuth **(99** parameters, I705 data, 17.2 observations/parameter).

Table 17 (Continued)

L.



 $\mathcal{A}$ 

 $\hat{\boldsymbol{\beta}}$ 



CL19 0.445(1) 0.086(2) 0.302(5)

CL20 0.466(1) 0.274(1) 0 0.463(1) 0.280(2) 0.465(2) 0.274(3)

 $0.443(1)$   $0.083(1)$   $0.305(4)$  $0.441(1)$   $0.084(2)$   $0.302(5)$ 

Table 17 (Continued)



 $\label{eq:2} \frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} \left(\frac{1}{\sqrt{2}}\right)^{2}$ 



 $\label{eq:2.1} \mathcal{L}(\mathcal{L}^{\text{max}}_{\mathcal{L}}(\mathcal{L}^{\text{max}}_{\mathcal{L}})) \leq \mathcal{L}(\mathcal{L}^{\text{max}}_{\mathcal{L}}(\mathcal{L}^{\text{max}}_{\mathcal{L}}))$ 

**Service** 

 $\label{eq:2.1} \frac{1}{2} \int_{\mathbb{R}^3} \frac{1}{\sqrt{2}} \, \frac{1}{\sqrt{2}} \,$ 

 $\mathcal{A}$  $\ddot{\phantom{a}}$   $\frac{1}{\sqrt{2}}$ 

 $\label{eq:2.1} \frac{1}{2} \sum_{i=1}^n \frac{$ 

 $\hat{\mathcal{A}}$ 

The nonabismuth cations found in the compounds  $B1_{12}C1_{14}$ and Bi<sub>10</sub>Hf<sub>3</sub>Cl<sub>18</sub> are compared in Table 19 and Figure 15. The appreciable distortion of the ion found in the former compound was attributed to the effects of close neighboring chlorine Interactions (17). There are nine chlorine-bismuth contacts at distances between 3.22 and 3.37 Å in  $B1_{12}Cl_{14}$  whereas the shortest contact found for the less distorted cation in the hafnium salt was  $3.374$  Å. In addition, the high symmetry of the space group in which  $Bi_{10}Hf_3Cl_{18}$  crystallizes requires that even these longer interactions are symmetrically arranged about the cation. The average value for each of the two ions of symmetry Independent distances in the idealized structure are very similar (see Table 19).

The attempted refinements have not improved the agreement between the model and the data. Comparison of the computed errors in the positional parameters and the standard deviations of the calculated distances indicate that the original refinement gave the best agreement between the data and the model. The use of an independent set of data to refine the model and the agreement between the results of the refinements corroborate the structure. The largest difference between the bond distances computed by the different refinements have the largest standard deviations and are associated with the light atoms; *i.e.*, the distance between BI7 and CL14 computed from the film data decreases from 2.73(3) to 2.66(5) computed from the counter data. The absence of any apparent similarity



Table 19. Comparison of the Bi<sup>st</sup> cations in  $BI_{12}CI_{14}$  and  $BI_{10}HF_3CI_{18}$  (in A)

 $a_{\text{Distances}}$  reported in (17) based on isotropic refinement,

^Standard deviation in, parentheses refer to the least significant figure.

'The number of times the distance repeats within the cation.

Figure 15. Comparison of the Bi, cations found in  $Bi_{12}Cl_{14}$  (on the right) and  $Bi_1oHf_3Cl_1s$  (on the left). Polyhedra shape have been idealized. (Data for  $Bi_{12}Cl_{14}$  taken from  $(17)$ ).

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between the final Fourier difference functions calculated with the two sets of data suggests that the peaks remaining in the maps reflect errors associated with the data and not a systematic error in the structure.

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## THE LEWIS ACID STABILIZATION EFFECT

Lewis acid complexatlon has been used to vary the nature of the anion present In the solid compounds and to probe the effect that this change has on the stability of the polynuclear bismuth cations which are formed. The general Increase In the stability of the lower oxidation states of metals to dlsproportionation with increasing anion size has been attributed to the decrease in the difference between the lattice energies of the oxidized and reduced compounds (69). This rationalization is not valid in cases where the two states of the metal have different anions. When mononegative complex anions are formed, the additional stability is determined by the competition between the loss in lattice energy and the heat of association released on formation of the complex anion. It is evident from Table 20 that the addition of the first halide to a Lewis acid is energetically favorable but that successive additions are less favorable. In cases where a multicharged anion is formed, the complexatlon energy is considerably reduced but the higher ionic charge contributes to the lattice energy of the product. It is the interplay of these factors which aid in the formation of the unusual compounds discussed in this thesis.

The average oxidation state of bismuth in the compound  $(Bi^+) (Bi_s^+)(HfCl_s^-)$  is the same as in the  $Bi_s^3$ <sup>+</sup> ion found in the tetrachloroaluminate(III) system. An explanation is possible as to why the  $Bi_5^{3+}$  ions are not present without





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 $\label{eq:2.1} \begin{split} \frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} &\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} &\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} &\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} &\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} &\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} &\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} &\frac{$ 

consideration of the relative stabilization energies of the Lewis acids. Salts of dinegatlve anions will have half as many anions per unit charge of the cation as will salts containing mononegatlve anions. The ionic packing in  $(Bi_{\circ}^{5})$ ,  $(BiCl_{5}^{2-})$ ,  $(Bi_{\circ}Cl_{6}^{2-})$  and  $(Bi_{\circ}^{+})(Bi_{\circ}^{5+})(HfCl_{6}^{2-})$ , reveal a single sheath of anions about the cations which prevents cation-cation contacts. Many ionic structures have anionanion contacts, however, only those compounds which have extensive covalent interactions between the cations are known to have cation-cation contacts. In the hexachlorohafnate(IV) salt the large  $B15^+$  cations and the HfCl<sup>2</sup><sup>-</sup> anions appear to control the packing. The  $Bi^+$  ions fill the small holes in the lattice and, by their presence, require more anions to achieve charge neutrality of the compound. These small ions pack more efficiently than would a second polynuclear cation and therefore do not contribute to the problem of cationic interactions. In contrast bismuth subchloride has large  $Big^-$  anions which occupy more space per unit charge than the HfCl<sup>2</sup><sup>-</sup> ions; thus, sufficient anions are available to prevent cation-cation contacts without the presence of the bismuth(I) ion. In the hypothetical compound  $(Bi_5^{\,3^+})_2(HfCl_6^{2^-})_3$ the ratio of polynuclear cations to anions is 2:3, probably too low to prevent the cations from having extensive interactions.

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## ACKNOWLEDGEMENTS

The author wishes to thank Dr. John Corbett for his Initial suggestion for this work and for his guidance throughout Its completion.

He also wishes to acknowledge: helpful discussions on crystallography with Drs. Clardy, Hackert, and Jacobson; assistance with the collection of intensity data from James Benson; guidance in the use of the Faraday balance from Dr, Harold Schobert; discussions and cooperation regarding this research with past and present members of Physical and Inorganic Chemistry Group IXj and the assistance of Sue Musselman with the preparation of the manuscript.

With all this help and assistance this work would not have been completed without the encouragement and support of his wife. Sue, who delayed her own further education so that he might achieve his goals. The help and sacrifice are deeply appreciated.

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